

$$E \gg m^2$$

$$H = \frac{p^2}{2m}$$

$$P \rightarrow \hat{P}$$

$$H \rightarrow i\hbar \frac{\partial}{\partial t}$$

$$i\hbar \frac{\partial}{\partial t} |\Psi\rangle = \frac{p^2}{2m} |\Psi\rangle$$

$$(i\hbar \frac{\partial}{\partial t}) |\Psi\rangle = (4|\Psi\rangle)$$

Relatividad: $E = \sqrt{p^2 + m^2 c^4} \rightarrow H$

$$i\hbar \frac{\partial}{\partial t} |\Psi\rangle = (\sqrt{p^2 + m^2 c^4})^{1/2} |\Psi\rangle$$

$$(\sqrt{p^2 + m^2 c^4})^2 = m^2 \left(\frac{p^2}{m^2 c^2} + 1 \right)^{1/2} = m^2 \left(1 + \frac{p^2}{m^2 c^2} - \frac{1}{2} \frac{p^4}{m^4 c^4} + \dots \right)$$

problemas de causalidad convergencia

Relatividad especial: x, t isotrop. linealiz., ec. Schrödinger no relativa $\frac{p^2}{2m} \rightarrow \vec{p}^2$

1) Klein Gordon: $E = \sqrt{p^2 + m^2 c^4}$, $H = \vec{p}^2 + m^2 c^2$

$$(i\hbar \frac{\partial}{\partial t})^2 \rightarrow H^2$$

$$i\hbar \frac{\partial}{\partial t} |\Psi\rangle = (\vec{p}^2 + m^2 c^2) |\Psi\rangle = \left(-\frac{\hbar^2 \vec{\nabla}^2}{m^2} + m^2 c^2 \right) |\Psi\rangle$$

$$\lambda_c = \frac{\hbar}{mc}$$

una sola componente? Describe $s=0$ (planos, ...)

2) Dirac: $s=1/2$

$$\vec{p}^2 + m^2 c^2 = (e_x p_x + e_y p_y + e_z p_z)^2 + p_m^2 c^2 = (c \vec{p} + p_m c^2)^2$$

$$\langle \vec{p}_1 \vec{p}_2 \vec{p}_3 \rangle \times \vec{m}^2 c^2 = \langle (e_x p_1^x + e_y p_1^y + e_z p_1^z)^2 + p_m^2 c^2 \rangle + \langle (e_x p_2^x + e_y p_2^y + e_z p_2^z)^2 + p_m^2 c^2 \rangle + \langle (e_x p_3^x + e_y p_3^y + e_z p_3^z)^2 + p_m^2 c^2 \rangle$$

$$e_x^2 = e_y^2 = e_z^2 = 1 \quad \vec{p}^2 = 1$$

$$e_x e_y + e_y e_z + e_z e_x = 0 \quad \{e_x, e_y\} = 0 \quad \{e_x, e_z\} = 0 \quad \{e_y, e_z\} = 0$$

$$\{e_x, p\} = 0$$

Soluciones a partir de matrices $a|x\rangle$!!

$$H = c \vec{a} \cdot \vec{p} + p_m c^2$$

a, p matrizes hermitianas, vértices propios ± 1

$$a_{xy} = -a_{yx} \quad a_{yz} = -a_{zy} \quad a_{xz} = -a_{zx}$$

$$\text{Tr}(a_{xy} a_{xy}^\dagger) = \text{Tr}(-a_{xy})$$

$$\text{Tr}(a_{xy}) = -\text{Tr}(a_{xy})$$

\vec{a}, \vec{p} no son únicas $\vec{a} \rightarrow S \vec{a} S^{-1}$, S unitaria también válido.

$$\vec{p} \rightarrow S \vec{p} S^{-1}$$

Representaciones \vec{a}, \vec{p} :

$$\vec{a} = \begin{pmatrix} 0 & \vec{\sigma} \\ \vec{\sigma} & 0 \end{pmatrix} \quad \vec{p} = \begin{pmatrix} 0 & 0 \\ 0 & -i \vec{\sigma} \end{pmatrix} \quad \text{dim } \mathcal{H} = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}^{2 \times 2}$$

$$\vec{a}_x = \begin{pmatrix} 0 & \sigma_x \\ \sigma_x & 0 \end{pmatrix} \quad \vec{a}_{y,z} = \begin{pmatrix} 0 & \sigma_y \\ \sigma_y & 0 \end{pmatrix}$$

$$\vec{\sigma} = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad \vec{\sigma}^2 = 1 \quad \vec{\sigma}_x^2 = 1 \quad \vec{\sigma}_y^2 = 1 \quad \vec{\sigma}_z^2 = 1$$

función vectorial relativista: $i\hbar \frac{\partial}{\partial t} |\Psi\rangle = (\vec{a} \cdot \vec{p} + p_m c^2) |\Psi\rangle$

es Dirac $|\Psi\rangle = \begin{pmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \\ \psi_4 \end{pmatrix}$!!

$|\Psi\rangle$: espino, 4 pisos

$$f_{\alpha_1} : i\hbar \frac{\partial}{\partial t} \psi_{\alpha_1} = \underbrace{(c \vec{a} \cdot \vec{p} + p_m c^2)}_{\vec{p}^2 + m^2 c^2} \psi_{\alpha_1}$$

4 funciones complejas $\begin{pmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \\ \psi_4 \end{pmatrix} \xrightarrow{\vec{p}^2 + m^2 c^2 = 0} \begin{pmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \\ \psi_4 \end{pmatrix} \xrightarrow{\vec{p}^2 + m^2 c^2 = 0} \begin{pmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \\ \psi_4 \end{pmatrix} \xrightarrow{\vec{p}^2 + m^2 c^2 = 0} \begin{pmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \\ \psi_4 \end{pmatrix}$

$$\vec{p}^2 + m^2 c^2 = 0 \quad \vec{p} = c \vec{\sigma} \vec{\omega} \quad \text{entre bloques } \psi_1, \psi_2$$

es relati. Dirac: $\begin{pmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \\ \psi_4 \end{pmatrix} \rightarrow \begin{pmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \\ \psi_4 \end{pmatrix}, \quad \vec{p} = c \vec{\sigma} \vec{\omega}$

$e^+ \rightarrow "e_L"$

$\vec{p} \rightarrow \vec{p}^+$

\vec{p}

\vec{p}^+

\vec{p}

Interacciones: CEM

$$(\vec{A}, \phi) \quad , \quad H_{\text{class}} = \left((\vec{p} - \frac{q}{c} \vec{A})^2 c^2 + m^2 c^4 \right)^{\frac{1}{2}} + q\phi$$

$\vec{p} \rightarrow \vec{p} - \frac{q}{c} \vec{A}$ acoplamiento mínimo

$$i\hbar \frac{\partial \Psi}{\partial t} = \left(c \vec{\alpha} \cdot (\vec{p} - \frac{q}{c} \vec{A}) + \beta mc^2 + q\phi \right) \Psi$$

Interacción !! $\vec{\alpha} \cdot \vec{A}$

$$\Psi(k, t) = \Psi e^{-iEt/\hbar}$$

$$\Rightarrow E\Psi = (c \vec{\alpha} \cdot \vec{\pi} + \beta mc^2 + q\phi)\Psi \quad , \quad \vec{\pi} = \vec{p} - \frac{q}{c} \vec{A}$$

$$\vec{\alpha} = \begin{pmatrix} 0 & \vec{\sigma} \\ \vec{p} & 0 \end{pmatrix} \quad \Psi = \begin{pmatrix} \chi \\ \phi \end{pmatrix} \quad , \quad \chi, \phi \text{ son biespinoros} \quad \chi = \begin{pmatrix} \chi_1 \\ \chi_2 \end{pmatrix} \quad \phi = \begin{pmatrix} \phi_1 \\ \phi_2 \end{pmatrix}$$

Usar operador $\vec{\alpha}$, $\vec{\beta}$:

$$\begin{pmatrix} E - mc^2 - q\phi & -c \vec{\sigma} \cdot \vec{\pi} \\ -c \vec{\sigma} \cdot \vec{\pi} & E + mc^2 - q\phi \end{pmatrix} \begin{pmatrix} \chi \\ \phi \end{pmatrix} = 0 \hat{=} \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

$$(E - mc^2 - q\phi)\chi - c \vec{\sigma} \cdot \vec{\pi} \phi = 0$$

$$(E + mc^2 - q\phi)\phi - c \vec{\sigma} \cdot \vec{\pi} \chi = 0$$

$$\rightarrow \phi = \frac{c \vec{\sigma} \cdot \vec{\pi}}{E + mc^2 - q\phi} \chi$$